

Problem Set 7 Solutions

1. (a) Expand $x(\epsilon)$ as a perturbation series in ϵ :

$$x(\epsilon) = x_0 + x_1\epsilon + x_2\epsilon^2 + O(\epsilon^3) \quad (1)$$

Substitute into our equation and consider the equation at each order:

$$O(\epsilon^0) : \quad x_0^3 - 1 = 0$$

$$O(\epsilon^1) : \quad 3x_0^2x_1 - x_0 = 0 \Rightarrow x_1 = \frac{1}{3x_0}$$

$$O(\epsilon^2) : \quad 3(x_0^2x_2 + x_0x_1^2) - x_1 = 0 \Rightarrow x_2 = 0$$

The solutions to for x_0 are $x_0=1, -\frac{1}{2} \pm \frac{\sqrt{3}}{2}i$. Let's consider only the real root for now. We have found:

$$x(\epsilon) = 1 + \frac{1}{3}\epsilon + O(\epsilon^3) \quad (2)$$

For $\epsilon=0.001$, MATLAB gives $x(\epsilon)=1.00033333332099$ while our approximation gives $x(\epsilon)\approx 1.00033333333333$. These solutions differ by approximately 1.23×10^{-11} , confirming that our solution is valid at least to $O(\epsilon^3)$.

- (b) Expand $x(\epsilon)$ as a perturbation series in ϵ :

$$x(\epsilon) = x_0 + x_1\epsilon + x_2\epsilon^2 + O(\epsilon^3) \quad (3)$$

Substitute into our equation and consider the equation at each order:

$$O(\epsilon^0) : \quad x_0^3 - x_0 = 0$$

$$O(\epsilon^1) : \quad 3x_0^2x_1 + x_0^2 - x_1 = 0 \Rightarrow x_1 = \frac{x_0^2}{1 - 3x_0^2}$$

$$O(\epsilon^2) : \quad 3(x_0^2x_2 + x_0x_1^2) + 2x_0x_1 - x_2 = 0 \Rightarrow x_2 = \frac{x_0x_1(2 + 3x_1)}{1 - 3x_0^2}$$

The solutions for x_0 are $x_0=-1, 0, 1$. So near the roots we have:

$$x(\epsilon) = -1 - \frac{1}{2}\epsilon - \frac{1}{8}\epsilon^2 + O(\epsilon^3) \quad (4)$$

$$x(\epsilon) = 0 + O(\epsilon^3) \quad (5)$$

$$x(\epsilon) = 1 - \frac{1}{2}\epsilon + \frac{1}{8}\epsilon^2 + O(\epsilon^3) \quad (6)$$

Here is a table showing $x(\epsilon)$ as calculated with MATLAB compared to our expansion.

x_0	$x_0 + \epsilon x_1$	$x_0 + \epsilon x_1 + \epsilon^2 x_2$	$x(\epsilon)$
-1	-1.00050000000000	-1.00050012500000	-1.00050012499999
0	0	0	0
1	0.99950000000000	0.99950012500000	0.99950012499999

2. Write the system as a 2D dynamical system:

$$\dot{x} = v \quad (7)$$

$$\dot{v} = a - x + \mu(1 - x^2)v \quad (8)$$

The fixed point of this system is $(x,v)=(a,0)$. The Jacobian is:

$$J = \begin{pmatrix} 0 & 1 \\ -1 - 2\mu xv & \mu(1 - x^2) \end{pmatrix}$$

Evaluated at the fixed point:

$$J = \begin{pmatrix} 0 & 1 \\ -1 & \mu(1 - a^2) \end{pmatrix}$$

So we have:

$$\lambda_{\pm} = \frac{\mu(1 - a^2)}{2} \pm \sqrt{\frac{\mu^2(1 - a^2)^2}{4} - 1} \quad (9)$$

To have a Hopf bifurcation we need $\mu(1 - a^2)=0$ and $\frac{\mu^2(1 - a^2)^2}{4} - 1 < 0$. Notice that the first condition implies the second condition. So we only need $\mu(1 - a^2)=0$. This means that Hopf bifurcations occur on the curves $a=\pm 1$ and $\mu=0$. I checked this numerically and found it to be true.

3. The Jacobian is:

$$J = \begin{pmatrix} \mu + y^2 & 2xy - 1 \\ 1 - 2x & \mu \end{pmatrix}$$

Evaluate at the origin:

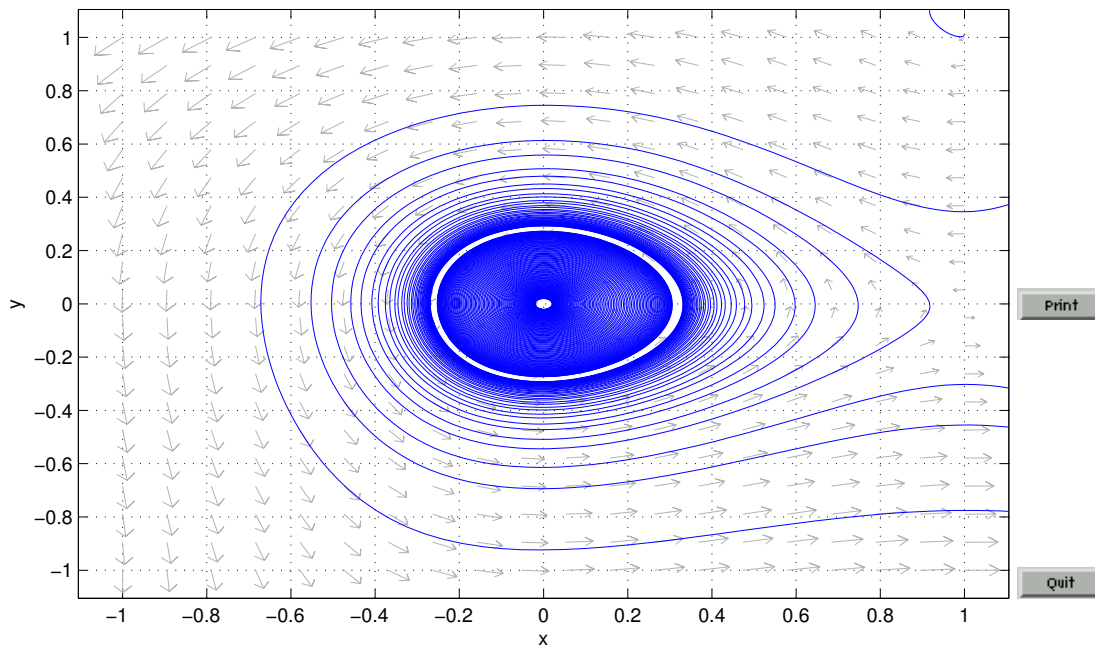
$$J = \begin{pmatrix} \mu & -1 \\ 1 & \mu \end{pmatrix}$$

So the eigenvalues are $\lambda_{\pm} = \mu \pm i$ and are pure imaginary when $\mu=0$.

4. There is an unstable limit cycle around the origin and a stable spiral at the origin for $\mu < 0$:

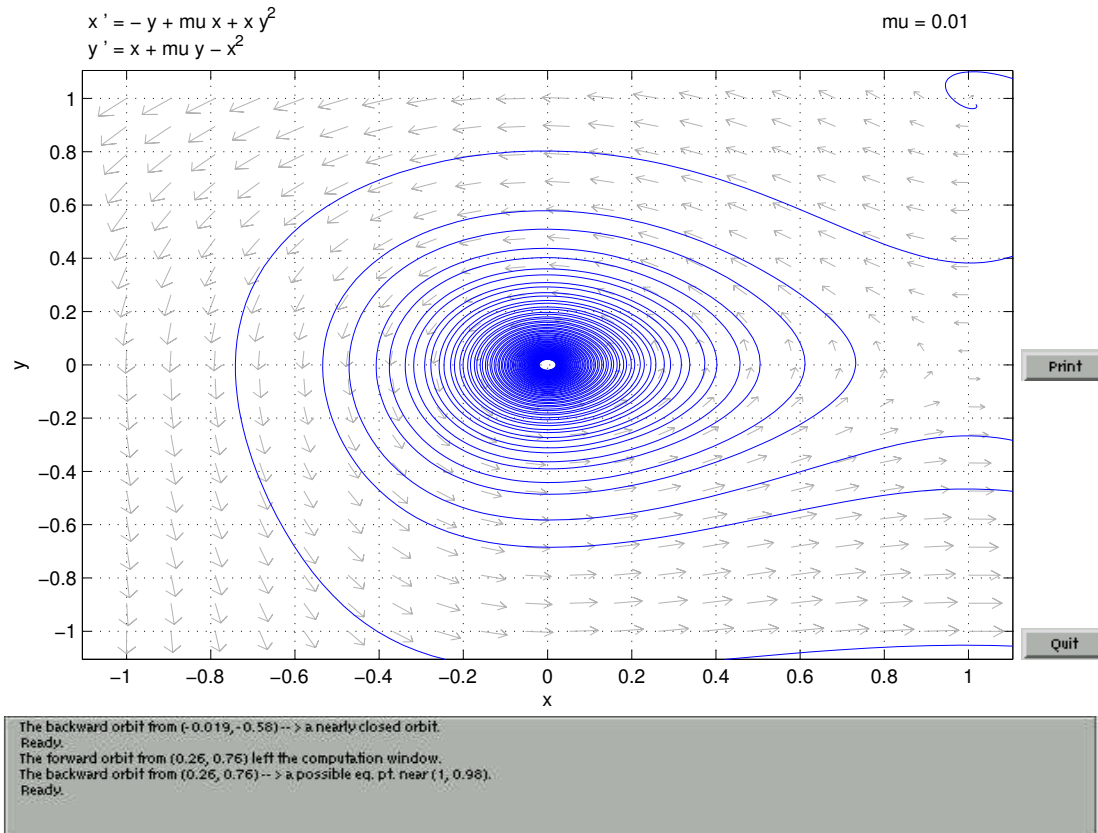
$$\begin{aligned}x' &= -y + \mu x + x y^2 \\y' &= x + \mu y - x^2\end{aligned}$$

$\mu = -0.01$



The backward orbit from (0.036, -0.067) --> a nearly closed orbit.
Ready.
The forward orbit from (0.29, 0.69) left the computation window.
The backward orbit from (0.29, 0.69) --> a possible eq. pt. near (1, 1).
Ready.

The origin is a unstable spiral and no limit cycle for $\mu > 0$. So a subcritical Hopf bifurcation occurs at $\mu = 0$.



5. (a) Use $x=r \cos(\theta)$ and $y=r \sin(\theta)$. The system can be rewritten:

$$\dot{r} = \mu r + r^2 \cos^2(\theta) \sin(\theta) [r \sin(\theta) - 1] \quad (10)$$

$$\dot{\theta} = 1 - r \cos(\theta) [\cos^2(\theta) + r^2 \sin^3(\theta)] \quad (11)$$

(b) Consider the average of these equations over θ . This will give us a qualitative idea of the behavior. We expect this method to capture the behavior best for $r \ll 1$ when θ is less important for the dynamics.

$$\dot{r} = \mu r + \frac{1}{8} r^3 \quad (12)$$

$$\dot{\theta} = 1 \quad (13)$$

(c) These equations suggest that an unstable limit cycle exists for negative μ with a radius of approximately $\sqrt{-8\mu}$ (consider the 1D dynamics in r) and the origin is unstable for positive μ . There is a subcritical Hopf bifurcation. Although not rigorous, this justifies the numerical results given above.